

## Decomposition of Mueller matrices in pure optical media

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### Abstract

The mathematical model of Mueller matrices is able to represent polarimetric properties of every material samples. In this paper an algebraic operation is performed to decompose a Mueller matrix  $M$  into the corresponding matrices of the pure optical media embedded into the complex material sample modeled by  $M$ .

**Keywords:** Polarized light, Mueller matrix.

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## 1 Introduction

In optical polarimetry, the state of polarization of a light beam is represented by the “Stokes vector” with four real elements arranged in a column vector. When a light beam interacts with a material medium, the Stokes vector  $\mathbf{s}$  that characterizes the incident light is transformed by the  $4 \times 4$  real Mueller matrix  $M$  that corresponds to the medium. The emerging light beam is characterized by another Stokes vector  $\mathbf{s}'$  given by the product  $\mathbf{s}' = M\mathbf{s}$ .

The mathematical structure of  $M$  depends on the complexity of the optical medium, so that we can distinguish between “pure Mueller matrices” and “non-pure Mueller matrices”. Some previous papers [1-5] deal with this property and its mathematical formulation.

The structure of Mueller matrices is studied by means of a transformation of  $M$  into a “coherency matrix”  $H$  given by [6]

$$H = \frac{1}{2} \begin{bmatrix} m_{00} + m_{01}+ & m_{02} + m_{12}+ & m_{20} + m_{21}- & m_{22} + m_{33}+ \\ m_{10} + m_{11} & i(m_{03} + m_{13}) & i(m_{30} + m_{31}) & i(m_{23} - m_{32}) \\ m_{02} + m_{12}- & m_{00} - m_{01}+ & m_{22} - m_{33}- & m_{20} - m_{21}- \\ i(m_{03} + m_{13}) & m_{10} - m_{11} & i(m_{23} + m_{32}) & i(m_{30} - m_{31}) \\ m_{20} + m_{21}+ & m_{22} - m_{33}+ & m_{00} + m_{01}- & m_{02} - m_{12}+ \\ i(m_{31} + m_{31}) & i(m_{23} + m_{32}) & m_{10} - m_{11} & i(m_{03} - m_{13}) \\ m_{22} + m_{33}- & m_{20} - m_{21}+ & m_{02} - m_{12}- & m_{00} - m_{01} \\ i(m_{23} - m_{32}) & i(m_{30} - m_{31}) & i(m_{03} - m_{13}) & -m_{10} + m_{11} \end{bmatrix}. \quad (1)$$

This expression indicates that there exists a simple linear relation between  $M$  and  $H$ , so that we can analyze the problem in terms of coherency matrices, which have a simpler mathematical characterization.

The degree and indices of purity of the material sample are given through the eigenvalues of  $H$ , so that  $H$  can be decomposed into a sum of one to four pure coherency matrices (i.e. the optical media can be represented by a combination of one to four pure material elements). Each pure coherency matrix contains a unique non-null eigenvalue, whereas non-pure coherency matrices contain 2, 3 or 4 non-null eigenvalues.

In this paper we deal with the mathematical resolution of a physical problem that appears frequently in polarimetry: Once obtained a measurement of the coherency matrix corresponding to the whole complex media, we want to “subtract” the action of a pure component that we know (or suspect) is present into the complex medium under measurement. We have developed a mathematical procedure to perform a proper subtraction and obtain the coherency matrix corresponding to the complex system regardless the effect of the known component.

## 2 Algebraic decomposition procedure

Given two positive semidefinite hermitic matrices  $H$  and  $A$  with dimension  $n$  such that  $\text{rang}A=1$ ,  $\text{rang}H=r$ ,  $0 < r \leq n$ , the following problem is stated:

“Is there exist a positive real number  $\alpha$  such that  $\text{rang}(H - \alpha A) = r - 1$ ?”

This question can be completely answered in two steps: regular  $H$  case, and general  $H$  case.

### Regular $H$ case

If  $\text{rang}H = n$ , then  $\alpha \neq 0$  is a necessary condition for  $\det(H - \alpha A) = 0$ . Therefore the problem stated before is equivalent to the problem of eigenvalues

$$\det\left(\frac{1}{\alpha}I - H^{-1}A\right) = 0. \quad (2)$$

It is easy to prove that only one eigenvalue  $1/\alpha > 0$  exists, whereas 0 is the other eigenvalue with multiplicity degree  $n - 1$ .

The solution to this problem also could be solved using the classical procedure for simultaneous diagonalization of two quadratic forms [7], extending it to the case of two hermitic forms with coordinate matrices  $H$  (positive definite) and  $A$  (positive semidefinite) in a base  $B = \{\mathbf{u}_1, \mathbf{u}_2, \dots, \mathbf{u}_n\}$ . The procedure consists of building an orthonormal base  $E = \{\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n\}$  by means of the Gram-Schmidt method applied to the hermitic product

$$P_H(\mathbf{z}, \mathbf{w}) = \mathbf{z}^T H \overline{\mathbf{w}}, \quad \mathbf{z}, \mathbf{w} \in \mathcal{C}^n. \quad (3)$$

The hermitic form associated with the matrix  $H$  expressed in the base  $E$  is

$$P_H(\mathbf{z}, \mathbf{z}) = \mathbf{z}^T \mathbf{z}, \quad (4)$$

and the hermitic form given by the matrix  $A$  is

$$P_A(\mathbf{z}, \mathbf{z}) = \mathbf{z}^T \begin{bmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_n \end{bmatrix} \mathbf{z}, \quad (5)$$

where  $\lambda_i$ ,  $i = 1, \dots, n$  are the eigenvalues of the hermitic matrix  $\overline{C}^T \overline{A} C$  and  $C$  is the matrix of the basis change from the base  $B$  to the base  $E$ . The equation  $\text{rang}(H - \alpha A) = n - 1$  formerly stated is solved computing  $\alpha$  such that  $1 - \alpha \lambda_i = 0$  for some  $i = 1, \dots, n$ , with  $\lambda_i \neq 0$ . Taking into account that  $\overline{C}^T \overline{A} C$  is a hermitic matrix, we have that  $\lambda_i$  are positive, and then  $\alpha > 0$ . The uniqueness of the solution  $\alpha$  is justified by the hypothesis  $\text{rang} A = 1$ .

### General $H$ case

From the simultaneous diagonalization procedure above described, we can approach the general case of  $\text{rang} H = r$ , where  $0 < r < n$ . If  $\text{rang} H = r < n$ , is not always possible to make the subtraction  $H - \alpha A$  with  $\text{rang}(H - \alpha A) = r - 1$ . On the light of this consideration, we see that now the goal is to characterize the existence of solution  $\alpha$  by means of the most efficient test.

To obtain the test proposed below, we need some suitable notation.

The matrix  $A$ , with  $\text{rang} A = 1$ , is denoted by  $A = (a_1 \mathbf{t}, a_2 \mathbf{t}, \dots, a_n \mathbf{t})^T$  where  $\mathbf{t}$  is a nonzero row vector.

The vector  $\mathbf{a} = (a_1, a_2, \dots, a_n)^T$  is named *the proportionality vector of the matrix  $A$* .

If  $L = [l_{kj}]$  is a order  $n$  regular matrix such that  $LH = (\mathbf{v}_1, \mathbf{v}_2, \dots, \mathbf{v}_r, 0, \dots, 0)^T$ , with  $\mathbf{v}_i$  row vectors, then:

$$L(H - \alpha A) = (\mathbf{v}_1 - \alpha \lambda_1 \mathbf{t}, \dots, \mathbf{v}_r - \alpha \lambda_r \mathbf{t}, -\alpha \lambda_{r+1} \mathbf{t}, \dots, -\alpha \lambda_n \mathbf{t})^T, \quad (6)$$

where  $\lambda_k = \sum_{j=1}^n l_{kj} a_j$ ,  $k = 1, \dots, n$ .

Denoting  $V_l = \text{span}\{\text{columns of } A\}$ ,  $V_r = \text{span}\{\text{rows of } H\}$  it is shown that  $\lambda_i = 0$  ( $i = r + 1, \dots, n$ ), is a necessary and sufficient condition to  $V_l \subset V_r$ , and there exists  $\alpha$  such that  $\text{rang}(H - \alpha A) = r - 1$ .

On the other hand, if  $\lambda_i \neq 0$  for some  $i = 1, \dots, n$ , then  $\mathcal{C}^{r+1} = V_l \oplus V_r$ , and we have that the subtraction  $H - \alpha A$  is not possible.

A similar test can be stated regarding the product  $H\bar{L}^T = (\mathbf{s}_1, \dots, \mathbf{s}_r, 0, \dots, 0)$  where  $\mathbf{s}_j$  ( $j = 1, \dots, n$ ), are column vectors. Taking the matrix  $A$  made by columns, both tests are equivalents.

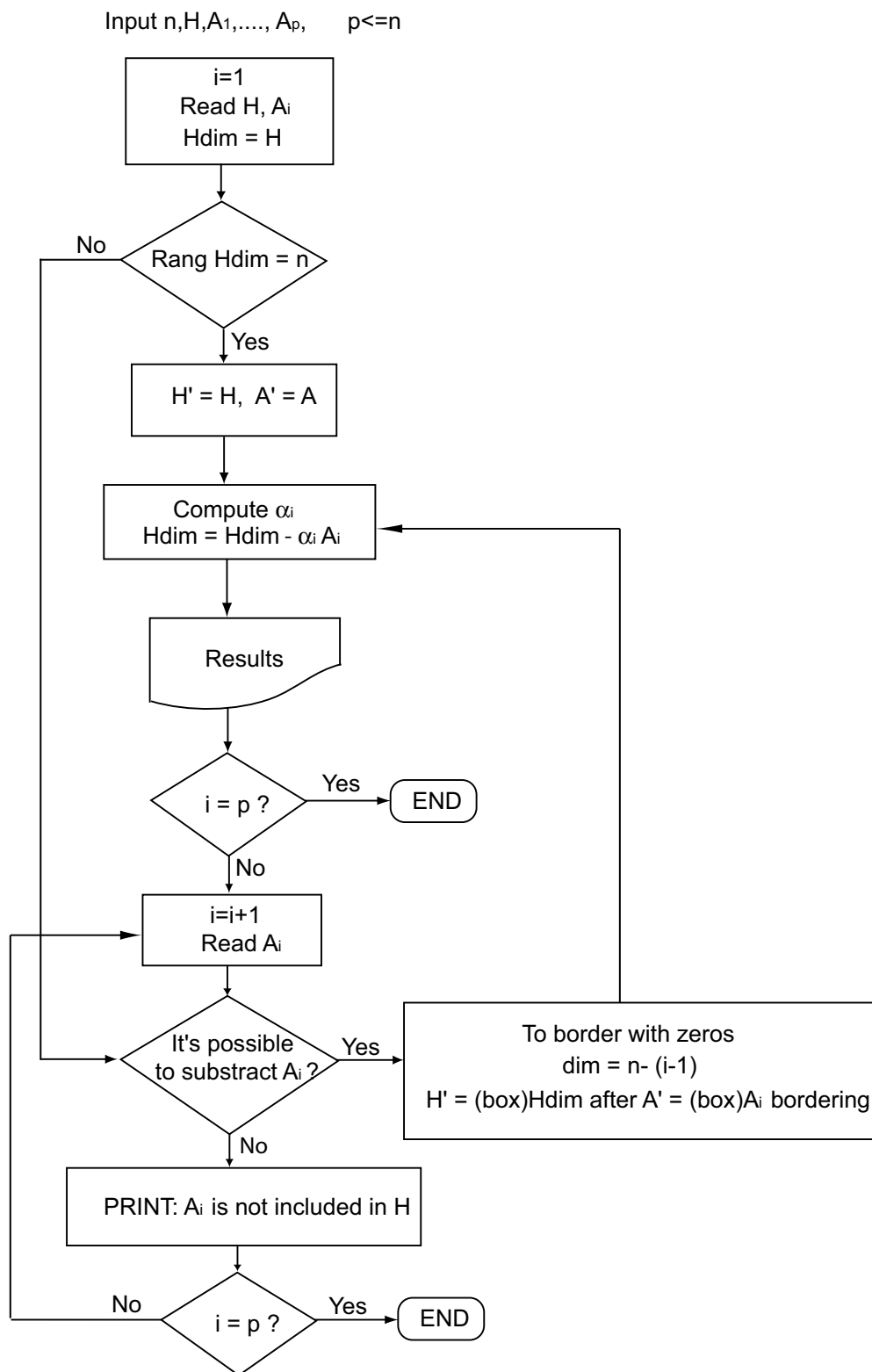
If the test performed with  $H$  and  $A$  has been successful, then the transformation  $L(H - \alpha A)\bar{L}^T$  leads, for every  $\alpha$ , to a matrix bordered with zeros in the last  $n - r$  rows and columns:

$$LH\bar{L}^T = \begin{bmatrix} \begin{bmatrix} H' \\ 0 \end{bmatrix} & 0 \\ 0 & 0 \end{bmatrix}, \quad LA\bar{L}^T = \begin{bmatrix} \begin{bmatrix} A' \\ 0 \end{bmatrix} & 0 \\ 0 & 0 \end{bmatrix}. \quad (7)$$

Now, considering the previously solved regular case, and taking  $H = H'$ ,  $A = A'$  and  $n = r$ , the general problem is solved and we can finally state:

“Given  $n \times n$  positive semidefinite hermitic complex matrices  $H$  and  $A$  such that  $\text{rang}A = 1$ ,  $\text{rang}H = r$ ,  $0 < r \leq n$ , there exists a unique real  $\alpha > 0$  such that  $\text{rang}(H - \alpha A) = r - 1$ ”.

### 3 Iterative scheme



## 4 An application example

In order to clarify the method, we present here an example of application of the procedure described in the previous sections.

We consider the matrix

$$H = \frac{1}{320} \begin{bmatrix} 53 & 15 & 15 & 61 \\ 15 & 53 & 5 - 48i & 15 \\ 15 & 5 + 48i & 53 & 15 \\ 61 & 15 & 15 & 53 \end{bmatrix}, \quad (8)$$

that corresponds to a non-pure material, with  $\text{rang}H = 3$ , and the matrix

$$P = \frac{1}{32} \begin{bmatrix} 9 & 3 & 3 & 9 \\ 3 & 1 & 1 & 3 \\ 3 & 1 & 1 & 3 \\ 9 & 3 & 3 & 9 \end{bmatrix}, \quad (9)$$

that corresponds to a pure polarizer [8]. In this case, the test is successful and the subtraction is possible with  $\alpha_1 = 1/2$ .

This value of  $\alpha_1$  represents the proportion of the material represented by  $P$  into the complex material represented by  $H$ . Then, the rest is

$$H - \frac{1}{2}P = \frac{1}{40} \begin{bmatrix} 1 & 0 & 0 & 2 \\ 0 & 6 & -6i & 0 \\ 0 & 6i & 6 & 0 \\ 2 & 0 & 0 & 4 \end{bmatrix}, \quad (10)$$

and it can be tested if the pure material represented by

$$P_1 = \frac{1}{8} \begin{bmatrix} 1 & 0 & 0 & 2 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 2 & 0 & 0 & 4 \end{bmatrix}, \quad (11)$$

is a component of the rest (10). Now, we can verify the test again and we obtain  $\alpha_2 = 1/5$ . This value means the concentration of polarizer  $P_1$  [8] into the material  $H$ .

Following the iterative procedure, the difference

$$H - \frac{1}{2}P - \frac{1}{5}P_1 = \frac{3}{20} \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & -i & 0 \\ 0 & i & 1 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}, \quad (12)$$

is a new matrix, which range is 1, that constitutes a retarder [8] and represents the remaining rest material:

$$R = \frac{3}{20} \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & -i & 0 \\ 0 & i & 1 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}. \quad (13)$$

It is easily tested that others optically pure materials are not present into the complex material represented by matrix  $H$ . For example, if the test is made for a retarder given by the matrix [8]

$$R_l = \frac{1}{2} \begin{bmatrix} 1 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 1 \end{bmatrix}, \quad (14)$$

the result of the test is negative and the subtraction is not possible.

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